

Introduction

In 2014, Ambainis [1] formalized a very natural physical problem: Given local Hamiltonian H and observable A , how difficult is it to *simulate* the measurement A on the ground space of H ? Formally:

APX-SIM [1]: Given k -local Hamiltonian H and l -local observable A , and $a, b, \delta \in \mathbb{R}$ such that $b - a \geq \frac{1}{\text{poly}(n)}$, $\delta \geq \frac{1}{\text{poly}(n)}$, for n the number of qubits H acts on, decide whether:

- YES: there exists a ground state $|\psi\rangle$ of H such that $\langle \psi | A | \psi \rangle \leq a$;
- NO: for all $|\psi\rangle$ s.t. $\langle \psi | H | \psi \rangle \leq \lambda_{\min}(H) + \delta$, it holds that $\langle \psi | A | \psi \rangle \geq b$.

Ambainis [1] showed that APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete for $O(\log n)$ -local H and A , where:

$\text{P}^{\text{QMA}[\log]}$: the set of problems decidable in polynomial time given $O(\log n)$ queries to a QMA oracle.

Improvements:

- G. and Y. showed [2] showed APX-SIM remains $\text{P}^{\text{QMA}[\log]}$ -complete even for **5-local Hamiltonian H** and **1-local measurement A** .
- [2] also showed $\text{P}^{\text{QMA}[\log]}$ is only “slightly harder” than QMA, in that $\text{P}^{\text{QMA}[\log]} \subseteq \text{PP}$.

Motivating question:

Does simulating measurements on ground spaces (APX-SIM) remain $\text{P}^{\text{QMA}[\log]}$ -complete for more physically motivated local Hamiltonians?

Main Results

We answer our motivating question positively. This is done via [Result 3](#), which first requires [Results 1 and 2](#).

Result 1: Parallel vs. adaptive queries. We show that $O(\log n)$ *adaptive* queries to a StoqMA or QMA oracle is equivalent to $\text{poly}(n)$ *parallel* queries to the oracle. Formally:

$$\text{P}^{\text{StoqMA}[\log]} = \text{P}^{\text{StoqMA}} \quad \text{and} \quad \text{P}^{\text{QMA}[\log]} = \text{P}^{\text{QMA}}.$$

Result 2: Complexity of \forall -APX-SIM under simulations. We show that the complexity of a *seemingly* easier problem, \forall -APX-SIM (see proof techniques), is preserved under “simulations” (in the sense of [3]).

Combined with known simulation results [3], this yields several complexity classifications for our original problem APX-SIM: It is in P, or is $\text{P}^{\text{NP}[\log]}$, $\text{P}^{\text{StoqMA}[\log]}$, or $\text{P}^{\text{QMA}[\log]}$ -complete for several Hamiltonian families.

Result 3: Complexity of APX-SIM for physical Hamiltonians.

Leveraging [Result 1](#), we show APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete for spatially-sparse H .

Combining with [Result 2](#), APX-SIM remains $\text{P}^{\text{QMA}[\log]}$ -complete for any Hamiltonian family which can efficiently simulate spatially-sparse H .

Punchline: APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete on physically motivated models like the Heisenberg anti-ferromagnetic interaction on a 2D lattice.

Proof techniques

Our approach proceeds in two high-level steps:

1. Give $\text{P}^{\text{QMA}[\log]}$ an equivalent characterization in terms of *polynomially many parallel* queries, i.e. $\text{P}^{\text{||QMA}}$, which eases the analysis of using Ambainis’s [1] query Hamiltonian construction ([Result 1](#)).
2. We wish to apply the “simulation” framework of [3] to show that APX-SIM is $\text{P}^{\text{||QMA}}$ -complete on physically motivated H . Four substeps:
 - a) Introduce intermediary problem, \forall -APX-SIM (see def. below).
 - b) Show that simulation preserves the complexity of \forall -APX-SIM ([Result 2](#)).
 - c) Show that \forall -APX-SIM is $\text{P}^{\text{||QMA}}$ -complete for spatially sparse H .
 - d) Apply existing simulation results [3] to obtain [Result 3](#), i.e. that APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete for various physically motivated models.

Notes:

- **Simulation** [3]: H_0 is a simulation of Hamiltonian H if there exists an efficiently computable local isometry $V = \otimes_i V_i$ that maps eigenvectors and eigenvalues of H to those of H_0 with “sufficiently small errors”.
- \forall -APX-SIM: Defined as APX-SIM but with “ \forall low-energy states $|\psi\rangle$ ” in the YES case. This problem is more robust than APX-SIM to perturbations in the ground space.

$\text{P}^{\text{QMA}[\log]} = \text{P}^{\text{||QMA}}$

Classical Theorem [4]: $\text{P}^{\text{NP}[\log]} = \text{P}^{\text{||NP}}$.

Forward direction $\text{P}^{\text{C}[\log]} \subseteq \text{P}^{\text{||C}}$ was shown for all classes C.

Generate all $2^{\log(n)} = \text{poly}(n)$ possible queries and pass to $\text{P}^{\text{||C}}$ machine.

Reverse direction $\text{P}^{\text{NP}[\log]} \supseteq \text{P}^{\text{||NP}}$ was shown using NP oracle to binary search the number of parallel queries which are YES-instances.

This technique fails in the quantum case! The P machine may make *invalid* queries, i.e. which violate promise gap, ex. k -LH with $\lambda_{\min} \in (a, b)$.

Theorem: $\text{P}^{\text{QMA}[\log]} = \text{P}^{\text{||QMA}}$

Proof: Forward direction as above. To show the reverse direction (\supseteq), we leverage a *hardness* result. Given that APX-SIM $\in \text{P}^{\text{QMA}[\log]}$, we prove that APX-SIM is $\text{P}^{\text{||QMA}}$ -hard.

We do so by adapting the “query Hamiltonian” constructions of [1], [2]:

$$H'_{\text{queries}} = \sum_{i=1}^m (2\epsilon|0\rangle\langle 0|_{x_i} \otimes I + |1\rangle\langle 1|_{x_i} \otimes H_{\text{query}}^i)$$

Additional benefits:

- Here, we use the classical Cook-Levin construction rather than Kitaev’s circuit-to-Hamiltonian (as in [2]). This yields $O(1)$ promise gap.
- This indirect method can be seen as simplifying the original proof that APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete, by greatly easing analysis of H'_{queries} .

This technique works for any class C for which there exists a family of Hamiltonians for which k -LH is C-complete! ex. NP, StoqMA, or QMA

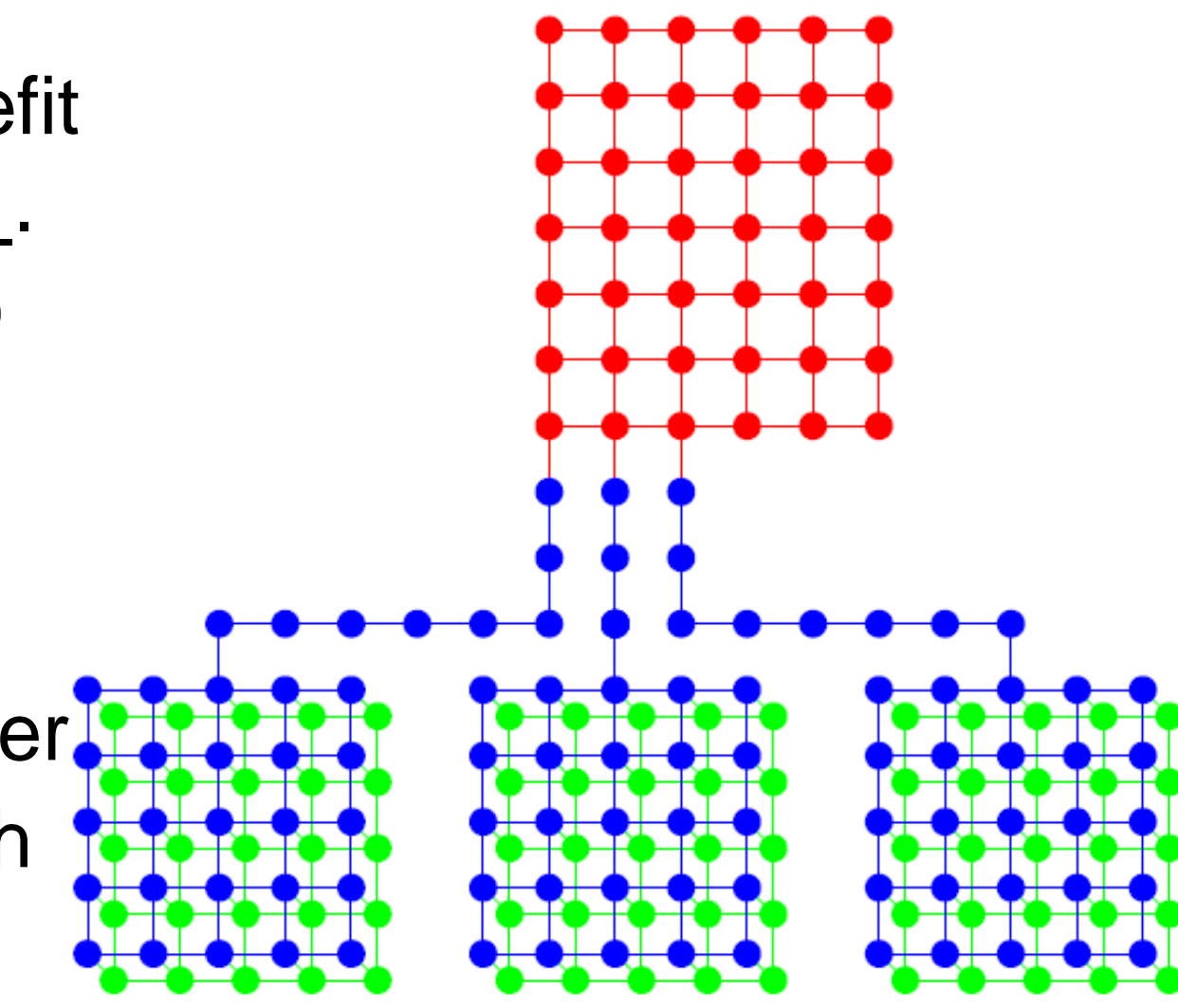
Complexity of APX-SIM for physical H

Theorem: APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete even for H restricted to a spatially-sparse interaction graph (in the sense of [5]).

Proof: We modify the hardness construction from [Result 1](#) so that it is spatially-sparse. There are 3 Hamiltonian terms,

$$H_{\text{final}} = H_{\text{Cook-Levin}} + H_{\text{queries}} + H_{\text{stab}}:$$

- $H_{\text{Cook-Levin}}$ is already spatially sparse (on a 2D lattice, in fact) – another benefit of our modifications in proving [Result 1](#).
- Since k -LH is QMA-hard for H on a 2D lattice [5], we assume all query instances H_{query}^i in H_{queries} are spatially-sparse.
- But, H_{queries} also has an answer register such that answer qubit x_i interacts with every qubit in H_{query}^i ; this is not spatially-sparse. We “spread out” the answer register to a grid of qubits overlaid on the query register, and introduce a stabilizing term H_{stab} .



Combined with [Result 2](#) and prior results for simulating spatially-sparse Hamiltonians, we obtain many corollaries. Here are a few:

Corollaries: APX-SIM is $\text{P}^{\text{QMA}[\log]}$ -complete even for H of XY interactions; for H of Heisenberg interactions; or for H on a 2D square lattice.

Discussion

Key takeaways:

1. The natural problem of simulating $O(1)$ -qubit measurements against ground states of physically motivated systems, such as the Heisenberg XY and antiferromagnetic interaction on a 2D lattice, is harder than QMA.
2. Working with *parallel* queries, as opposed to *adaptive* queries, simplifies circuit-to-Hamiltonian constructions.

Open questions:

- What other results for k -LH / QMA can be extended to $\text{P}^{\text{QMA}[\log]}$?
- Identify additional $\text{P}^{\text{QMA}[\log]}$ -complete problems and physical inputs.
- What happens if we switch the P base with other classes (classical or quantum) (ex. $\text{BQP}^{\text{QMA}[\log]}$)?
- What if we use other quantum oracles (ex. $\text{P}^{\text{BQP}[\log]}$)?

References

- [1] A. Ambainis, CCC 2014. [3] T. Cubitt, A. Montanaro, and S. Piddock, PNAS 2017.
 [2] J. Yirka and S. Gharibian, TQC 2017. [4] R. Beigel, TCS 1991.
 [5] R. Oliveira and B. Terhal, QIC 2008.

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